

**FY3403 Particle physics**

NTNU

**Problemset 7**

Institutt for fysikk

**SUGGESTED SOLUTION****Problem 1**

The traces are evaluated by using the rules derived in page 238-240, chapter 7 Griffiths. Using *e.g.* that the trace of an odd number of  $\gamma$ -matrices is zero, we arrive at:

$$\langle |\mathcal{M}|^2 \rangle = \frac{8g_e^4}{(p_1 - p_3)^4} [(p_1 p_2)(p_3 p_4) + (p_1 p_4)(p_2 p_3)] \quad (1)$$

in the high-energy limit where we set  $(m, M) \rightarrow 0$ . The remaining part is to evaluate the products of the 4-vectors in the CM frame. Assume that  $\theta$  is the scattering angle between  $p_1$  and  $p_3$ . Due to a net momentum of zero before and after the collision, we may write:

$$p_1 = (E_1/c, \mathbf{p}_1), p_2 = (E_2/c, -\mathbf{p}_1), p_3 = (E_3/c, \mathbf{p}_3), p_4 = (E_4/c, -\mathbf{p}_3). \quad (2)$$

However, since the masses are neglected and energy is conserved, we have  $E_j = E$  and hence  $|\mathbf{p}_1| = |\mathbf{p}_3|$ . Inserting this into the expression for the amplitude above, we find:

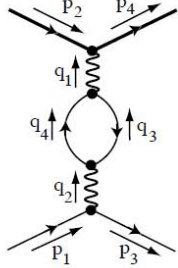
$$\langle |\mathcal{M}|^2 \rangle = \frac{8g_e^4}{(p_1 - p_3)^4} (E^4/c^4) [4 + (1 + \cos \theta)^2]. \quad (3)$$

Moreover, we have that  $(p_1 - p_3)^4 = (2\mathbf{p}^2 - 2\mathbf{p}^2 \cos \theta)^2 = 4\mathbf{p}^4 (1 - \cos \theta)^2$ . Now use that in the CM frame for two-body scattering (result derived previously in the lectures/book), we have:

$$\frac{d\sigma}{d\Omega} = (\hbar c / 8\pi)^2 \frac{S \langle |\mathcal{M}|^2 \rangle}{(E_1 + E_2)^2} \quad (4)$$

Here, the statistical factor is  $S = 1$ . Inserting our expression for the amplitude into the above equation, we find after some trigonometric manipulations that:

$$\frac{d\sigma}{d\Omega} = (\hbar c / 8\pi)^2 \frac{g_e^4}{2E^2} \frac{1 + \cos^4 \theta / 2}{\sin^4 \theta / 2} \quad (5)$$

**Problem 2**

Applying the Feynman rules,

$$\begin{aligned} & \int [\bar{u}(3) (ig_e \gamma^\mu) u(1)] \frac{-ig_{\mu\lambda}}{q_2^2} [\text{LOOP}] \frac{-ig_{\lambda\nu}}{q_1^2} [\bar{u}(4) (ig_e \gamma^\nu) u(2)] \\ & \times (2\pi)^4 \delta^4(p_1 - p_3 - q_2) (2\pi)^4 \delta^4(q_2 - q_3 - q_4) (2\pi)^4 \delta^4(q_3 + q_4 - q_1) \\ & \times (2\pi)^4 \delta^4(q_1 + p_2 - p_4) \frac{d^4 q_1}{(2\pi)^4} \frac{d^4 q_2}{(2\pi)^4} \frac{d^4 q_3}{(2\pi)^4} \frac{d^4 q_4}{(2\pi)^4} \end{aligned}$$

where “LOOP” stands for

$$-\text{Tr} \left[ (ig_e \gamma^\lambda) \frac{i(\not{q}_4 + mc)}{q_4^2 - m^2 c^2} (ig_e \gamma^\kappa) \frac{i(\not{q}_3 + mc)}{q_3^2 - m^2 c^2} \right]$$

The  $q_2$  integral, using  $\delta^4(p_1 - p_3 - q_2)$ , sends  $q_2 \rightarrow p_1 - p_3 \equiv q$  (for short). The  $q_1$  integral, using  $\delta^4(q_1 + p_2 - p_4)$ , sends  $q_1 \rightarrow p_4 - p_2$ , and the two remaining delta functions  $\Rightarrow q_3 + q_4 = p_1 - p_3 = p_4 - p_2$ , so  $q_1$  is also  $q$  (of course). The  $q_3$  integral, using  $\delta^4(q - q_3 - q_4)$ , sends  $q_3 \rightarrow q - q_4$ , and we erase the final delta function  $(2\pi)^4 \delta^4(p_1 - p_3 + p_2 - p_4)$ . There is still an integral over  $q_4$ , which (for simplicity) we rename  $k$ . Multiplying by  $i$ :  $\mathcal{M} =$

$$-\frac{ig_e^4}{q^4} [\bar{u}(3) \gamma^\mu u(1)] \left\{ \frac{d^4 k}{(2\pi)^4} \frac{\text{Tr} [\gamma_\mu (\not{k} + mc) \gamma_\nu (\not{q} - \not{k} + mc)]}{[k^2 - m^2 c^2][(q - k)^2 - m^2 c^2]} \right\} [\bar{u}(4) \gamma^\nu u(2)] \quad \checkmark$$

**NB!** The direction of the momentum  $\vec{q}_3$  above is opposite to what it should be: it must be directed along the arrow on the fermion line if we are to use the expression for the fermion propagator. As a result, it should not be slashed  $q$  minus slashed  $k$  inside the integrand, but slashed  $k$  minus slashed  $q$  in the nominator (for the denominator, the order doesn't matter due to the square).