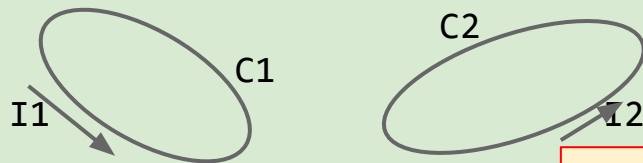


Inductance

Consider 2 coils of wire



The magnetic field through C1 is

$$\Phi_1 = L_1 I_1 + M_{12} I_2$$

and similarly, the flux through C2 is

$$\Phi_2 = L_2 I_2 + M_{21} I_1$$

This is important because

$$\frac{d\Phi_1}{dt} = L_1 \frac{dI_1}{dt} + M_{12} \frac{dI_2}{dt} = -(\text{emf})_1$$

When currents change, that will affect those and other currents by “inductance”.

L_i = self inductance
(how I_i affects itself)

M_{ij} = mutual inductance
(how I_j affects I_i)

$$B_1 \propto I_1$$

$$B_2 \propto I_2$$

Theorem $M_{12} = M_{21}$

Proof

$$\begin{aligned}\Phi_{12} &= M_{12} I_2 = \int_{S_1} \vec{B}_2 \cdot d\vec{a}_1 \\ &= \int_{S_1} (\nabla \times \vec{A}_2) \cdot d\vec{a}_1 \\ &= \oint_{C_1} \vec{A}_2 \cdot d\vec{\ell}_1\end{aligned}$$

Recall the vector potential,

$$\vec{A}_2(\vec{x}_1) = \oint_{C_2} \frac{\mu_0 I_2 d\vec{\ell}_2}{4\pi |\vec{x}_1 - \vec{x}_2|}$$

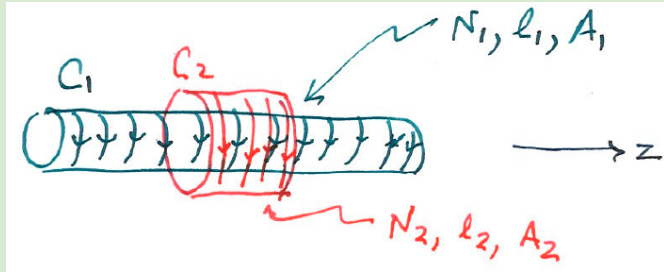
Thus,

$$M_{12} = \frac{\mu_0}{4\pi} \oint_{C_1} \oint_{C_2} \frac{d\vec{\ell}_1 \cdot d\vec{\ell}_2}{|\vec{x}_1 - \vec{x}_2|}$$

"Neumann's equation"

$$M_{12} = \frac{\mu_0}{4\pi} \oint_{C_1} \oint_{C_2} \frac{d\vec{l}_1 \cdot d\vec{l}_2}{|\vec{x}_1 - \vec{x}_2|} \quad \text{Q. E. D.}$$

Example. The mutual inductance between coaxial solenoids, Sol₁ and Sol₂, assuming these are densely wound solenoids and Sol₂ is outside Sol₁ ...



Assume $l_2 < l_1$ and $A_2 > A_1$.
Neglect end effects.

Case 1 Consider $\Phi_{21} = \int_{S_2} \vec{B}_1 \cdot d\vec{a}_2 = M_{21} I_1$

$$\vec{B}_1 = \mu_0 \frac{N_1}{l_1} I_1 \hat{k} \text{ in area } A_1$$

$$\Phi_{21} = \mu_0 \frac{N_1}{l_1} I_1 N_2 A_1 = M_{21} I_1$$

$$\therefore M_{21} = \mu_0 N_1 N_2 \frac{A_1}{l_1} \quad (1)$$

Case 2 Consider $\Phi_{12} = \int_{S_1} \vec{B}_2 \cdot d\vec{a}_1$

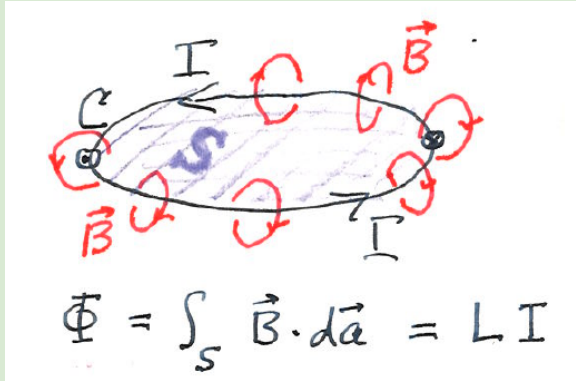
$$\vec{B}_2 = \mu_0 \frac{N_2}{l_2} I_2 \hat{k} \text{ in area } A_2 \text{ and length } l_2$$

$$\Phi_{12} = \left(\mu_0 \frac{N_2}{l_2} I_2 \right) \left(\underbrace{N_1 \frac{l_2}{l_1}}_{\substack{\text{number of turns of } C_1 \\ \text{in length } l_2}} \right) A_1 = M_{12} I_2$$

$$\therefore M_{12} = \mu_0 N_1 N_2 \frac{A_1}{l_1} \quad (2)$$

Note that $M_{12} = M_{21}$

Self Inductance (L)

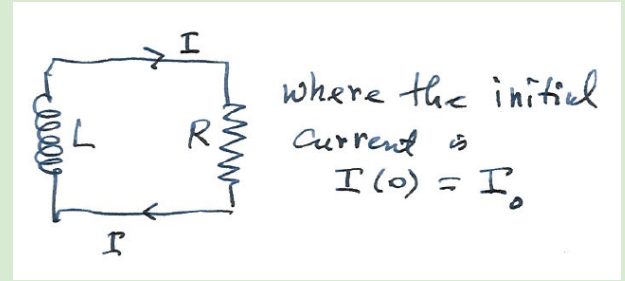


Example: a solenoid (= a wire densely wound around a long cylinder)

$$\vec{B} = \mu_0 \frac{N}{l} I \hat{k}$$
$$\Phi = \int_S \vec{B} \cdot d\vec{a} = \mu_0 \frac{N}{l} I NA$$
$$L = \mu_0 N^2 \frac{A}{l}$$

Units: $\text{Tm/A} \times \text{m}^2/\text{m} = \text{Tm}^2/\text{A} = \text{H}$ (henry)

LR Circuit



The emf around the inductor is

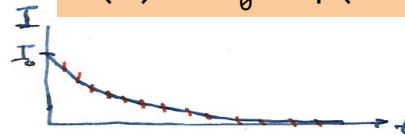
$$\text{emf} = - \frac{d\Phi}{dt} = - \frac{d}{dt}(LI)$$
$$= -L \frac{dI}{dt}$$

This emf drives current

$$I = \frac{\text{emf}}{R} \quad (\text{Ohm's law})$$

(just like a battery!). Thus

$$I = \text{emf} / R = - d/dt (LI) / R$$
$$dI / dt = - (R/L) I$$
$$I(t) = I_0 \exp(-Rt/L)$$



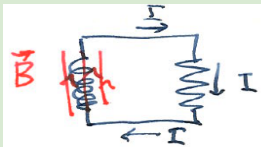
Energy density of a magnetic field

Apply energy conservation to the LR circuit. The power dissipated in resistance is $P = I^2 R$ (Joule's law).

Conservation of energy implies

$$\frac{dU}{dt} = -P$$

where U is the energy of the magnetic field in the inductor. So



$$\frac{dU}{dt} = -RI_0^2 e^{-2Rt/L}$$

$$P dt = dQ V$$
$$P = IV = I^2 R$$

$$U - U_0 = \int_0^t (-RI_0^2) e^{-2Rt/L} dt$$
$$= -RI_0^2 \left(\frac{-L}{2R}\right) (e^{-2Rt/L} - 1)$$
$$= \frac{1}{2} LI^2 - \frac{1}{2} LI_0^2$$

Evidently, $U = \frac{1}{2} LI^2 = \frac{\Phi^2}{2L}$.

Recall L and Φ for a solenoid....

$$U = \frac{1}{2} \left(\frac{\mu_0 N^2 A}{l} \right) (NBA)^2$$

$$U = \frac{1}{2} \frac{B^2}{\mu_0} A l = \frac{1}{2} \frac{B^2}{\mu_0} \times \text{Volume}$$

The energy density is $u = \frac{U}{V} = \frac{B^2}{2\mu_0}$.